RUSSIAN JOURNAL OF EARTH SCIENCES vol. 14, 3, 2014, doi: 10.2205/2014ES000543

# Weighted gravitational time series smoothing

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**Abstract.** This article continues the series of papers by the authors on the new universal DMA-smoothing of time series, originally intended for the analysis of geophysical time series obtained in the framework of discrete mathematical analysis (DMA), developed by GC RAS. We formulated the general concept of weighted DMA-smoothing, constructed and analyzed one of its variants.

This is the e-book version of the article, published in the Russian Journal of Earth Sciences (doi:10.2205/2014ES000543). It is generated from the original source file using LaTeX's **epub.cls** class.

# Introduction

Time series are the most effective way to represent dynamic data, the results of various kinds of monitoring (physical, geophysical, environmental et al.). Smoothing of time series is one of the most powerful tools to study them, because the smoothed time series often are simpler, more amenable to simple analysis and serve as an effective approximation of the original time series.

To date, the most known methods of smoothing are regression analysis, finite element method and splines, Fourier and wavelet smoothing and their generalizations based on expansions in orthogonal systems of functions and time-frequency analysis.

The proposed method expands the list of a new type of smoothing, able to operate successfully in the analysis of geophysical time series. In the article it was applied to real data from the Borok geomagnetic observatory of the INTERMAGNET network.

## Main Concept of DMA-Smoothing

Let TS[a, b] be a space of time series on a discrete segment [a, b] with nodes  $t_i = a + (i-1)h$ , i = 1, ..., n, h = (b - a)/(n - 1). The elements of TS[a, b] are denoted by x, y, z, .... If  $z \in TS[a, b]$ , then  $z_i = z(t_i)$ and the correspondence  $z \sim (z_i)|_1^n$  turns TS[a, b] into *n*-dimensional Euklidian space  $R^n = R^n(z_1, ..., z_n)$ .

We considered the task of constructing the smoothing  $\text{Sm } y \stackrel{\text{def}}{=} x \in \text{TS}[a, b]$  for the given time series  $y \in \text{TS}[a, b]$ . Its solution is a consistent formalization of the following conjunction:

"x -smoothing y"  $\equiv (x -$ approximation  $y) \land (x -$ smoothed time series)

(1)

Let us denote by z a variable series in TS[a, b]. The approximation in (1) is formalized by the quadratic residual functional  $Sc(z|y) = ||z-y||^2$ , and the smoothness – by the non-negative quadratic form CG(z), which we call the smoothness residual. It is based on a variant of the discrete continuity. The residual CG(z) is a measure of the deviation from this property for the time series z on the interval [a, b]. The less are the values Sc(z|y) and CG(z), the more reason we have to consider z the smoothing y. Consequently, we need a meaningful simultaneous minimization of the functionals Sc(z|y) and CG(z). To do this, we considered their nontrivial convex combination

$$\mathsf{Sm}_\lambda(z|y) = \lambda \mathsf{CG}(z) + (1-\lambda)\mathsf{Sc}(z|y), \lambda \in (0,1)$$

and a search of the smoothing x = Sm y for y is reduced to a minimization on TS[a, b] at a suitable  $\lambda \in (0, 1)$  of the functional  $\text{Sm}_{\lambda}(z|y)$ , i.e. to the solution of the *n*-dimensional linear system

$$x = \operatorname{Sm}_{\lambda} y \Leftrightarrow \operatorname{Grad} \operatorname{Sm}_{\lambda}(x|y) = 0$$

Thus we suggest the whole family  $\{\operatorname{Sm}_{\lambda}y, \lambda \in (0,1)\}$  as the smoothing  $\operatorname{Sm} y$ . Let us find the gradient  $\operatorname{Grad} \operatorname{Sm}_{\lambda}(x|y)$  in its clear form. For that first we introduce the non-negative on  $\operatorname{TS}[a, b]$  operator G, which generates the gravitational smoothness residual  $\operatorname{CG}(z) = (\operatorname{G} z, z)$  and then we transform  $\operatorname{Sm}_{\lambda}(z|y)$ :

$$\operatorname{Sm}_{\lambda}(z|y) = \lambda(\operatorname{G} z, z) + (1 - \lambda)||z - y||^2 =$$

$$\lambda(\mathsf{G}\,z,z) + (1-\lambda)(z-y,z-y) =$$

 $\lambda(\operatorname{G} z, z) + (1 - \lambda)(z, z) - 2(1 - \lambda)(z, y) + (1 - \lambda)||y||^2$ 

 $((\lambda \mathsf{G} + (1-\lambda)\mathsf{E})z, z) - 2(1-\lambda)(z, y) + (1-\lambda)||y||^2$ 

Hence E is a unit operator on TS[a, b]. Consequently, the minimization  $Sm_{\lambda}(z|y)$  is equal to the

functional minimization

$$\widetilde{\mathsf{Sm}}_{\lambda}(z|y) = \frac{1}{2}((\lambda\mathsf{G} + (1-\lambda)\mathsf{E})z, z) - ((1-\lambda)z, y)$$

For it the gradient in the point *z* is expressed through G like this [*Pshenichny and Danilin*, 1975]:

$$\operatorname{\mathsf{Grad}} \operatorname{\widetilde{\mathsf{Sm}}}_\lambda(z|y) = (\lambda \mathsf{G} + (1-\lambda)\mathsf{E})z - (1-\lambda)y$$

Thus,

$$x = \mathsf{Sm}_{\lambda}y \Leftrightarrow (\lambda\mathsf{G} + (1-\lambda)\mathsf{E})x = (1-\lambda)y$$

If  $\lambda \in (0, 1)$ , then the operator  $\lambda G + (1 - \lambda)E$  is positive on TS[*a*, *b*], and therefore reversible. Consequently, the smoothing  $x = Sm_{\lambda}y$  always exists, is uniquely defined for any  $y \in TS[a, b]$  and represents the result of applying the operator  $(\lambda G + (1 - \lambda)E)^{-1}$ , therefore

$$\mathsf{Sm}_\lambda = (1-\lambda)(\lambda\mathsf{G} + (1-\lambda)\mathsf{E})^{-1}$$

Let us consider an important variant of the residual CG(z).

#### **Gravitational Residual of Smoothness**

Let us briefly (according to the module [Agayan et al., 2010; Gvishiani et al., 2011]) consequently, from top

to bottom recall the construction of the main, but not the only one residual of the discrete smoothness CG(z): the gravitational smoothness residual CGr(z):

$$\mathsf{CGr}(z) = \sum_{s=0}^{n-1} \omega_s \mathsf{CGr}^s(z)$$

where  $\omega_s$  is the non-negative weights (parameters of smoothness). The residual CGr(z) was generated by the operator of gravitational smoothness Gr:

$$\operatorname{CGr}(z) = (\operatorname{Gr} z, z), \ \operatorname{Gr} = \sum_{s=0}^{n-1} \omega_s \operatorname{Gr}^s$$

where

$$Gr^s = D^{s^*}Gr[a, b^s]D^s$$

and  $D^s$  is the *s*-dimensional discrete differentiation operator: if *z* is the time series on the interval [a, b], then its *s*-dimensional derivative  $z^s = D^s z$  is time series on the interval  $[a, b^s]$ ,  $b^s = b - sh$  and in each node  $t_i \in [a, b^s]$ 

$$z^{s}(t_{i}) = \frac{\sum_{l=0}^{s} (-1)^{s-l} C_{s}^{l} z_{i+l}}{h^{s}}$$

The operator  $Gr[a, b^s]$  is called the operator of gravitation continuity residual on the interval  $[a, b^s]$  and is defined by the proximity measure  $\delta_{t_i}^s(t_j)$  in accordance with the following formula

$$Gr[a, b^{s}] = (A(\delta^{s}) - E)^{*}(A(\delta^{s}) - E)$$
$$A(\delta^{s}) = (a_{ij}^{s})$$
$$a_{ij}^{s} = \frac{\delta_{i}^{s}(j)}{\sum_{j=1}^{n-s}(\delta_{i}^{s}(j))}$$

One of the following constructions was used as the proximity measure:

$$\delta_i^s(j) = \delta_{t_i}^s(t_j) = \left(1 - \frac{|t_j - t_i|}{\max(|t_i - a|, |t_i - b_s|) + h}\right)^p$$

as shown in Figure 1, or



Figure 1. Global proximity measure.

$$\delta_i^s(j) = \delta_i^{s,r}(j) = \delta_{t_i}^{s,r}(t_j) =$$

$$\begin{cases} \left(1 - \frac{|t_j - t_i|}{r}\right)^p & \text{if } |t_j - t_i| < r, \\ 0 & \text{if } |t_j - t_i| \ge r. \end{cases}$$

for a certain  $r < b_s - a$ , as shown in Figure 2

#### Parameters of Gravitational Smoothness

Thus there are three free parameters of gravitational smoothing: the coupling multiplier  $\lambda \in [0, 1]$ , the system w of non-negative weights  $w_s$ , s = 0, ..., n-1 and



Figure 2. Local proximity measure.

the system  $\delta$  of proximity measures  $\delta^s$ , s = 0, ..., n-1, which depends on the internal parameter  $p \ge 0$ .

These parameters give the gravitational smoothing greater flexibility and expressiveness. The study allows some qualitative conclusions about their relationship to be made:

- 1. the increment  $\lambda$  at fixed w and  $\delta$  leads to increasing smoothness and reduces the scanning (Figure 3).
- 2. the increment p at fixed  $\lambda$  and w leads to improving scanning with the preservation of smoothness (Figure 4).
- 3. the influence of weights w at fixed  $\lambda$  and p on gravitational smoothness is more complicated and



**Figure 3.** Dependence on  $\lambda$  at fixed w and  $\delta$ .



**Figure 4.** Dependence on p at fixed  $\lambda$  and  $\delta$ .

depends on their relationship. In Figure 5 the simplest dependence on basic weights is shown.

4. our experience with the algorithm allowed the most appropriate values of the free parameters to be identified:  $\lambda \in [0.9, 0.999]$ ,  $p \in [2, 10]$ ,  $w = \{w_0 = 0, w_{1-4} = 1, w_s = 0, s > 4\}$  (Figure 6).

# Weighted DMA-Smoothing. Main Concept

Let us assume that each node  $t_i$  in the interval [a, b] has its non-trivial weight  $v_i \ge 0$ . Let us denote by **V** their total population  $\mathbf{V} = \{v_i|_1^n\}$ . This raises the question of the impact of **V** on the DMA-smoothing, which is the subject of this paragraph. The functional

$$\mathsf{Sm}_{\lambda}(x|y) = \lambda \mathsf{CG}(x) + (1-\lambda)\mathsf{Sc}(x|y)$$

is responsible for the smoothing, so the impact **V** occurs separately: through the smoothness residual CG(x), the scanning functional Sc(x|y) and the connecting multiplier  $\lambda$ .

In the future we shall also denote by V the diagonal operator with weights  $v_i$ . Further, the smoothing



Observatory BOX, Component Y (14-Nov-2008)



Observatory BOX, Component Y (14-Nov-2008)

elements, into which the weights **V** have (or have not) penetrated, will be respectively denoted by  $CG_V(x)$ ,  $Sc_V(x|y)$ ,  $\lambda_V$  (CG(x), Sc(x|y),  $\lambda$ ). Relationships between the smoothing Sm and the weights V in a general case can develop according to one of the following eight scenarios:

$$Sm_{\lambda}(x|y) = \lambda CG(x) + (1 - \lambda)Sc(x|y)$$

$$VSm_{\lambda}(x|y) = \lambda CG(x) + (1 - \lambda)Sc_{V}(x|y)$$

$$SmV_{\lambda}(x|y) = \lambda CG_{V}(x) + (1 - \lambda)Sc(x|y)$$

$$VSmV_{\lambda}(x|y) = \lambda CG_{V}(x) + (1 - \lambda)Sc_{V}(x|y)$$

$$Sm_{\lambda_{V}}(x|y) = \lambda_{V}CG(x) + (1 - \lambda_{V})Sc(x|y) \quad (2)$$

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$$\mathsf{SmV}_{\lambda_{\mathsf{V}}}(x|y) = \lambda_{\mathsf{V}}\mathsf{CG}_{\mathsf{V}}(x) + (1 - \lambda_{\mathsf{V}})\mathsf{Sc}(x|y)$$

$$\mathsf{VSmV}_{\lambda_{\mathsf{V}}}(x|y) = \lambda_{\mathsf{V}}\mathsf{CG}_{\mathsf{V}}(x) + (1 - \lambda_{\mathsf{V}})\mathsf{Sc}_{\mathsf{V}}(x|y)$$

The interest in the structures (2) is explained by the following circumstances: in general stochastic setting smoothing of time series can be considered as a possible option of ideal course of the process, which is described by the given time series. If the disturbances of process in time are expressed by weights, the weighted smoothing gives an idea of the impact of these disturbances on the process.

The simplest is the impact of **V** on the scanning Sc(x|y): it is natural to assume that

$$\mathsf{Sc}_V(x|y) = \sum_{i=1}^n v_i (x_i - y_i)^2$$

If the weights  $v_i$  are positive, then the operator **V** is reversible and the smoothing x = Smy (2) looks like this

$$\operatorname{Sm}_{\lambda} y = (1 - \lambda)(\lambda \mathsf{G} + (1 - \lambda)\mathsf{E})^{-1}y$$

$$VSm_{\lambda} y = (1 - \lambda)(\lambda G + (1 - \lambda)E)^{-1} \mathbf{V}y$$

$$SmV_{\lambda} y = (1 - \lambda)(\lambda G_{V} + (1 - \lambda)E)^{-1}y$$

$$VSmV_{\lambda} y = (1 - \lambda)(\lambda G_{V} + (1 - \lambda)E)^{-1} \mathbf{V}y$$

$$Sm_{\lambda_{V}} y = (1 - \lambda_{V})(\lambda_{V}G + (1 - \lambda_{V})E)^{-1}y$$

$$VSm_{\lambda_{V}} y = (1 - \lambda_{V})(\lambda_{V}G + (1 - \lambda_{V})E)^{-1}\mathbf{V}y$$

$$SmV_{\lambda_{V}} y = (1 - \lambda_{V})(\lambda_{V}G_{V} + (1 - \lambda_{V})E)^{-1}y$$

 $\mathsf{VSmV}_{\lambda_{\mathsf{V}}} y = (1 - \lambda_{\mathsf{V}})(\lambda_{\mathsf{V}}\mathsf{G}_{\mathsf{V}} + (1 - \lambda_{\mathsf{V}})\mathsf{E})^{-1}\mathbf{V}y$ 

To summarize, there are seven non-trivial scenarios of varying degrees of penetration depth of the weights  $\mathbf{V}$  of nodes in the DMA-smoothing Sm. The simplest is the second scenario  $VSm_{\lambda}$  ((2) and (3)) in which the weights of nodes penetrate the scanning only.

We conclude this work with the analysis of such smoothing in the gravitational case.

# Weighted Gravitational Smoothing VSm $_\lambda$

In all the smoothings in Figure 7–Figure 9 the nodes of the main record y were supposed to be equal, having a

unit weight. Let us break this equality three times by giving

- in the first case to the points on the interval [18:30– 18:50] the weight 10.0;
- in the second case to the points on the interval [19:50-20:10] the weight 0.1;
- in the third unifying case to the points on the interval [18:30–18:50] the weight 10.0, to the points on the interval [19:50–20:10] the weight 0.1.

In each case at the same parameters  $\lambda$ , w and p we constructed the smoothings  $Sm_{\lambda} y$  and  $VSm_{\lambda} y$ . In Figure 7–Figure 9 the smoothing  $Sm_{\lambda} y$  is shown in blue and outside the intervals [18:30-18:50] and [19:50–20:10] is concealed by the smoothing VSm $_{\lambda}$  y, shown in red. Our study has shown that the situation in Figure 7–Figure 9 typical for the correlation of  $Sm_{\lambda}y$ and VSm $_{\lambda}$  y and allows the following conclusion to be made: under the same conditions  $\lambda$ , w, p the smoothings  $Sm_{\lambda} y$  and  $VSm_{\lambda} y$  have the same smoothness, at that the construction  $VSm_{\lambda} y$  is clearly dependent on the weights V: the higher is the weight  $v_i$  in the node  $t_i$ , the more carefully the smoothing VSm<sub> $\lambda$ </sub> y scans the record y in the vicinity  $t_i$ .



Observatory BOX, Component Y (14-Nov-2008)







**Figure 9.** The points' weights in the period 18:30-18:50 are equal to 10.0, in the period 19:50-20:10 are equal to 0.1.

## References

- Agayan, S. M., A. D. Gvishiani, Sh. R. Bogoutdinov, A. I. Kagan (2010), Smoothing of time series by methods of discrete mathematical analysis, *Russian Journal of Earth Sciences*, 11, RE40001, doi:10.2205/2009ES000436.
- Gvishiani, A. D., S. M. Agayan, Sh. R. Bogoutdinov, A. I. Kagan (2011), Gravitational smoothing of time series, *Institute* of Mathematics and Mechanics of the Ural Branch of the Russian Academy of Sciences, 17, 2, 62–70.
- Pshenichny, B. N., Yu. M. Danilin (1975), Numerical Methods in Extremal Tasks, Nauka, Moscow, 319.